Analysis of an Important Class of Non-Markov Systems

Yashwant K. Malaiya, Member IEEE
State University of New York, Binghamton
Stephen Y.H. Su, Senior Member IEEE
State University of New York, Binghamton

Key Words—Transient failures, Reliability analysis, Non-Markov processes.

Reader Aids-

Purpose: Develop methods for analysis Special math needed for explanation: Probability Special math needed to use results: Same Results useful to: Fault-tolerant computing, Reliability theoreticians

Summary & Conclusions—Probabilistic modeling of many types of systems generally assumes Markov behavior. However, some important practical systems exhibit memory. For example, in digital computer systems, the probability of occurrence of a transient failure is related to the time period the system has been operating correctly. Analytic methods do not yet exist that allow accurate modeling of such systems for the purpose of reliability analysis and fault-tolerant design.

Methods are presented here to analyze an important class of non-Markov systems. In this class, the transition-probability-rate of an outward transition from a state is related to the duration the system has continuously been in that state. To analyze such systems, concept of memory profile has been introduced. Methods are first presented which enable computation of steady-state probabilities for both discrete-time and continuous-time processes with two states. These are then extended for general non-steady-state cases and also for systems with more than two states.

I. INTRODUCTION

A wide range of systems has been modeled by Markov processes. For discrete-state continuous-time Markov process, the transition-probability-rate outwards from a state is governed by a constant parameter. This implies that the continuous-time duration spent in a state is exponentially distributed [1]. The advantage of assuming a process to be Markov is that the mathematical analysis is tractable [2-5]. The arrival of transient (intermittent) faults in digital circuits has so far been analyzed assuming a Markov process. However, in a study of multiple-processor systems [6], the Weibull distribution is found to have a better match with experimental data. This implies that in such systems, the transition-probability rate from a state depends on the time the system has spent in that state. The methods presented here allow analysis of this important class of non-Markov systems.

The motivation in developing the methods was to analyze transient failures in digital systems more accurately. The methods are very general, and are not restricted to Weibull (which is more general than exponential) distribution. They can be used to compute the reliability and the availability of digital systems with (or without) redundancy.

They can also be used to design proper test experiments for digital systems with transient faults. Because of their generality, these methods can also be used for other systems which exhibit similar non-Markov behaviour.

To examine a system in which transition-probabilityrate depends on the duration the system has been in that state, requires considering the history of each state. We shall first consider, in section II, discrete-time systems to develop the necessary concepts; continuous-time systems are then examined in section III. The results for the continuous-time systems are next extended for non-steadystate and then finally for multi-state cases.

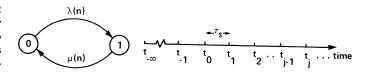


Fig. 1. Discrete time system with two states.

II. DISCRETE-TIME SYSTEMS

Consider a discrete-time system portrayed in figure 1.

Assumptions

- 1. It is a discrete-time 2-state system.
- 2. The parameter governing the transition out from a state depends on the time the system has spent in that state.
 - 3. The system is in steady-state.

The steady-state probability of the system being in either state will be found here.

Notation

r_s time-step

- $\lambda(n)$ Pr{system is in state 1 at instant t_j | (it was in state 0 at instant t_{j-1}) \cap (system was in state 0 for a period $n\tau_s$)}
- $\mu(n)$ Pr{system is in state 0 at instant t_j | (it was in state 1 at instant t_{j-1}) \cap (system was in state 1 for a period $n\tau_s$)}
- $P_x(j)$ Pr{system is in state x at instant t_j }; x = 0 or 1
- $\pi_x(i, j)$ Pr{(system is in state x at instant t_i) \cap (it entered state x $i \cdot \tau_s$ time ago)}; x = 0 or 1

The function $\pi_x(i, j)$, i = 1 to ∞ is called *memory-profile*, and describes the probabilistic history of state x, x = 0 or 1, at instant t_j . The first index i refers to the history (with the dimensions of time) and the second index j refers to the instant in time, when the memory profile is defined.

It was necessary to define $\pi_x(i, j)$ so that the effect of memory could be represented.

The sum of all the values of the memory-profile $\pi_x(i, j)$, i = 1 to ∞ is $p_x(t_i)$.

Proof: By definition we have—

 $p_x(j) = \Pr\{\text{system in the state } x \text{ at instant } t_i\}$

- = $\sum_{i=1}^{\infty} \Pr\{\text{system in state } x \text{ at instant } t_i | \text{it entered} \}$ state $x i \cdot \tau_s$ time ago}. Pr{it entered state $x i \cdot \tau_s$ time ago }
- = $\sum_{i=1}^{\infty} \Pr\{(\text{system is in state } x \text{ at instant } t_i) \cap (\text{it en-}$ tered state $x i \cdot \tau_s$ time ago)}

$$= \sum_{i=1}^{\infty} \pi_x(i,j). \qquad QED \quad (1)$$

In order to obtain expressions for $p_x(j)$, x = 0 or 1, all $\pi_{x}(i, j)$ have to be evaluated. We first obtain an expression for $\pi_x(i, j)$, $i \ge 2$ in terms of $\pi_x(1, j)$. Then we evaluate $\pi_x(1, j)$ i) itself.

Consider for $i \ge 2$:

 $\pi_{x}(i, j) = \Pr\{\text{system is in state 0 with history } i \cdot \tau_{s} \text{ at time } i$

- = $\Pr\{\text{system is in state 0 with history } (i-1) \cdot \tau_s \text{ at } t_j\}$
- Pr{a transition to state 1 at time t_i |history is $(i-1)\cdot\tau_s$

$$= \pi_0(i-1,j) - \lambda(i-1) \pi_0(i-1,j)$$

= $[1 - \lambda(i-1)] \pi_0(i-1,j)$

$$= [1 - \lambda(i-1)] \pi_0(i-1,j)$$
 (2)

The solution to (2) is:

$$\pi_0(i,j) = \pi_0(i,j) \prod_{k=1}^{i-1} (1 - \lambda(k)), i \ge 2$$
(3a)

$$\pi_1(i, j) = \pi_1(1, j) \prod_{k=1}^{i-1} (1 - \mu(k)), i \ge 2$$
 (3b)

Now an expression will be obtained for $\pi_x(1, j)$, x = 0 or 1. Consider,

 $\pi_0(1, j) = \Pr\{\text{system is in state 0 with history } \tau_s \text{ at time } t_i\}$

- = $\sum_{k=1}^{\infty} \Pr\{\text{a transition to state 0 at time } t_{j-1} | \text{(sys-}$ tem in state 1 at t_{j-1}) \cap (it entered state 1 at $k\tau_s$ time ago)} • Pr{(system in state 1 at t_{j-1}) \cap (it entered state 1 at $k\tau_s$ time ago)} $= \sum_{k=1}^{\infty} \mu(k) \pi_1(k, j-1).$

As in the steady-state case $\pi_1(k, j - 1) = \pi_1(k, j)$,

$$\pi_0(1,j) = \sum_{k=1}^{\infty} \mu(k) \, \pi_1(k,j),$$
 (4a)

$$\pi_1(1,j) = \sum_{k=1}^{\infty} \lambda(k) \, \pi_0 - (k,j).$$
 (4b)

We will next show that $\pi_0(1, j) = \pi_1(1, j)$. Using (3b) and (4a), we have,

$$\pi_{0}(1,j) = \left\{ \sum_{k=2}^{\infty} \mu(k) \cdot \pi_{1}(1,j) \prod_{k'=1}^{k-1} (1 - \mu(k')) \right\}$$

$$+ \mu(1) \pi_{1}(1,j)$$

$$= \pi_{1}(1,j) \left[\sum_{k=2}^{\infty} \mu(k) \prod_{k'=1}^{k-1} (1 - \mu(k')) + \mu(1) \right]$$

$$= \pi_{1}(1,j).$$
(5)

The last step follows from the fact that the term within the brackets is equal to one as shown in the following proof.

Proof: Pr{there will be a transition to state 0 in infinite time | system was in state 1 at time 0 with history τ_s

$$= \mu(1) + \mu(2)(1 - \mu(1)) + \mu(3)(1 - \mu(1))(1 - \mu(2)) + \dots$$
$$= \mu(1) + \sum_{k=1}^{\infty} \mu(k) \prod_{k=1}^{k-1} (1 - \mu)(k'))$$

However the l.h.s. is equal to one. Hence, the r.h.s is equal **OED** to one.

Since
$$p_0(j) + p_1(j) = 1$$
, using (1), (3), (5),

$$\pi_0(1,j)[\Lambda+M] = 1 \tag{6}$$

$$\Lambda = 1 + \sum_{i=2}^{\infty} \prod_{k=1}^{i-1} (1 - \lambda(k))$$

$$M = 1 + \sum_{i=2}^{\infty} \prod_{k=1}^{i-1} (1 - \mu(k))$$
(7a)
(7b)

$$M \equiv 1 + \sum_{i=2}^{\infty} \prod_{k=1}^{i-1} (1 - \mu(k))$$
 (7b)

From (6) and (5)

$$\pi_0(1, j) = \pi_1(1, j) = 1/(\Lambda + M)$$

and hence

$$p_0(j) = \Lambda/(\Lambda + M), p_1(j) = M/(\Lambda + M)$$
 (8)

Once $\lambda(k)$ and $\mu(k)$ are given as functions of k, the steady-state probabilities $p_0(j)$ and $p_1(j)$ can be calculated using (8).

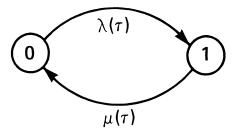


Fig. 2. Continuous time system with two states.

III. CONTINUOUS-TIME SYSTEMS

Continuous-time systems, represented in figure 2, can be similarly analyzed.

Assumptions

- 1. It is a continuous-time two-state system.
- 2. The parameter governing the transition out from a state depends on the time the system has spent in that state.
 - 3. The system is in steady-state.

The steady-state probability of the system being in either state will be found here.

Notation

 $\lambda(\tau)dt$ Pr{system is in state 1 at time $t \mid \text{(it was in state 0 at time } t - dt) \cap \text{(system was in state 0 for a period } \tau)}$

 $\mu(\tau)dt$ Pr{system is in state 0 at time $t \mid \text{(it was in state 1 at time } t - dt) \cap \text{(system was in state 1 for a period } \tau)}$

 $p_x(t)$ Pr{system is in state x at time t}; x = 0 or 1 $\pi_x(\tau, t)dt$ Pr{system is in state x at time t) \cap (it entered state x between τ to $\tau + d\tau$ time ago}

The memory profile $\pi_x(\tau, t)$, $\tau = 0$ to ∞ , is now continuous. For continuous cases, the $\lambda(\tau)$, $\mu(\tau)$, $\pi_0(\tau, t)$ all have dimensions of (time)⁻¹. Corresponding to (1), we now have—

$$p_x(t) = \int_0^\infty \pi_x(\tau, t) d\tau, x = 0, 1$$
 (9)

which means that $p_x(t)$ is the area under the memory-profile curve for time t.

For the steady-state case, we can get, corresponding to (2) and (3a)

$$\frac{\partial \pi_0(\tau, t)}{\partial \tau} = -\lambda(\tau)\pi_0(\tau, t), \tau > 0$$
 (10)

$$\pi_0(\tau, t) = \pi_0(0, t) \exp[-\int_0^{\tau} \lambda(\tau') d\tau']$$
 (11a)

Similarly,

$$\pi_1(\tau, t) = \pi_1(0, t) \exp[-\int_0^{\tau} \mu(\tau') d\tau']$$
 (11b)

Define-

$$\Lambda \equiv \int_0^\infty \exp[-\int_0^\tau \lambda(\tau')d\tau']d\tau \tag{12a}$$

$$M \equiv \int_0^\infty \exp[-\int_0^\tau \mu(\tau')d\tau']d\tau. \tag{12b}$$

Using (9) and (11),

$$p_0(t) = \pi_0(0, t)\Lambda \tag{13a}$$

$$p_1(t) = \pi_1(0, t)M \tag{13b}$$

Corresponding to (4), we now have—

$$\pi_0(0, t) = \int_0^\infty \pi_1(\tau', t) \mu(\tau') d\tau'$$
 (14a)

$$\pi_1(0, t) = \int_0^\infty \pi_0(\tau', t) \lambda(\tau') d\tau'. \tag{14b}$$

The memory-profiles can be eliminated as in the previous section. The steady-state probabilities, corresponding to (8) are:

$$p_0(t) = \Lambda/(\Lambda + M), p_1(t) = M/(\Lambda + M). \tag{15}$$

When $\lambda(\tau)$ and $\mu(\tau)$ are given as a function of τ , (15) can be used to compute the steady-state probabilities. In the derivation above, $\lambda(\tau)$ and $\mu(\tau)$ are any general functions of τ . Two important special cases are examined below.

When the duration of staying in a state has an exponential distribution, then both λ and μ are constant with respect to history τ , and the system is described by a Markov process. In this case

$$p_0(t) = \mu/(\lambda + \mu) \text{ and } P_1(t) = \lambda/(\lambda + \mu)$$
 (16)

which are well known results. For the Weibull distribution assume that:

$$\lambda(\tau) = \lambda' \alpha (\lambda' \tau)^{\alpha - 1}$$

$$\mu(\tau) = \mu' \beta (\mu' \tau)^{\beta - 1}$$

where α , β , λ' , μ' are constant parameters. In this case;

$$p_0(t) = \mu' \Gamma\left(\frac{1+\alpha}{\alpha}\right) / \left[\mu' \Gamma\left(\frac{1+\alpha}{\alpha}\right) + \lambda' \Gamma\left(\frac{1+\beta}{\beta}\right)\right]$$
 (17a)

$$p_1(t) = 1 - p_0(t) (17b)$$

These can be easily verified by considering the average time spent in each state. Only for 2-state system is the average time spent in each state directly related to the steady-state probabilities.

Equations (16) or (17) are very useful for analyzing fault-tolerance with transient (intermittent) failures.

IV. EXTENSION TO NON-STEADY-STATE CASE

Assumptions

- 1. It is a 2-state continuous time system.
- 2. The parameter governing the transition out from a state depends on the time the system has spent in that state.
 - 3. The system need not be in steady-state.

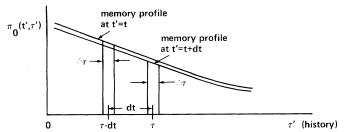


Fig. 3. Continuous memory profile for general case.

Consider the memory profiles at times t and t + dt, as shown in figure 3, these memory profiles might not be identical in general, as steady state is no longer assumed. Both $\Delta \tau$ and dt are infinitesimal, though not necessarily equal. Consider a small band of probability around history

 τ -dt of width $\Delta \tau$ at a time t. The area enclosed in this band is equal to the probability that the system is in state 0 at time t, and it entered state 0 between (t-dt)- $(\Delta \tau)/2$ and $(\tau-dt)+(\Delta \tau)/2$ time ago. If an outgoing transition from state 0 were not allowed, then at time t+dt, a band of width $\Delta \tau$ around τ would represent the same probability. However, since transitions from state 0 are allowed, the band at τ , bounded by the memory profile at t+dt, would be shorter than the one at τ -dt bounded by the memory profile at t. The difference would be related to the probability rate of outgoing transitions from state 0.

$$[\pi_0(\tau, t + dt) - \pi_0(\tau - dt, t)] \Delta \tau = -\pi_0(\tau - dt, t) \Delta \tau \lambda(\tau - dt),$$
 for $\tau > 0$

$$[\pi_0(\tau, t + dt) - \pi_0(\tau, t) + \pi_0(\tau, t) - \pi_0(\tau - dt, t)]$$

= $-\pi_0(\tau - dt, t)\lambda(\tau - dt)$, for $\tau > 0$

$$\pi_0(0, t)dt = dt \int_0^\infty \mu(\tau)\pi_1(\tau, t)d\tau$$
, for $\tau = 0$.

Following the usual calculus procedures, we have

$$\frac{\partial \pi_0(\tau, t)}{\partial t} = -\frac{\partial \pi_0(\tau, t)}{\partial \tau} - \pi_0(\tau, t)\lambda(\tau), \text{ for } \tau > 0 \text{ (18a)}$$

$$\frac{\partial \pi_0(0, t)}{\partial t} = \int_0^\infty \mu(\tau) \frac{\partial \pi_1(\tau, t)}{\partial t} d\tau, \text{ for } \tau = 0.$$
 (18b)

The steady state case can be obtained from (18a) by taking

$$\frac{\partial \pi_0(\tau, t)}{\partial t} = 0.$$
i.e.
$$\frac{\partial \pi_0(\tau, t)}{\partial \tau} = -\pi_0(\tau, t)\lambda(\tau), \text{ for } \tau > 0$$

which is identical to (10). For steady-state in (18b), a similar procedure yields

$$\frac{\partial \pi_1(\tau, t)}{\partial t} = 0.$$

The system probabilities are completely defined by (18) for each state, including the initial conditions. There is some redundancy in information if an additional equation is added to indicate that the total probability of being in either state is unity.

V. EXTENSION TO MULTIPLE STATES

Equation (18) can be extended to more than two states.

Assumptions

- 1. It is a multiple-state continuous time system.
- 2. The parameters governing the transition out from a state depends on the time the system has spent in that state.
 - 3. The system need not be in steady-state.

Let a transition from state k to state 1 be governed by the parameter $\lambda_{k1}(\tau)$. Then corresponding to (18), we have

$$\frac{\partial \pi_i(\tau, t)}{\partial t} = -\frac{\partial \pi_i(\tau, t)}{\partial \tau} - \pi_i(\tau, t) \sum_j \lambda_{ij}(\tau), \text{ for } \tau > 0$$
(19a)

$$\pi_i(0, t) = \sum_{j \neq i} \int_0^\infty \lambda_{ji}(\tau) \pi_j(\tau, t) d\tau, \text{ for } \tau = 0.$$
 (19b)

For the steady-state case,

$$\frac{\partial \pi_i(\tau, t)}{\partial \tau} \doteq -\pi_i(\tau, t) \sum_j \lambda_{ij}(\tau), \text{ for } \tau > 0$$

$$\pi_i(\tau, t) = \pi_i(0, t) \exp[-\int_0^\tau \sum_j \lambda_{ij}(\tau')d\tau'], \text{ for all } i \quad (20a)$$

Using this in (19b) yields

$$\pi_{i}(0, t) = \sum_{j \neq i} \pi_{j}(0, t) \int_{0}^{\infty} \lambda_{ji}(\tau)$$

$$\exp\left[-\int_{0}^{\tau} \sum_{k} \lambda_{jk}(\tau')d\tau'\right]d\tau, \text{ for all } i$$
(20b)

Equation (20b) is a simple system of linear algebraic equations in $\pi_i(0, t)$, for all *i*. Once all $\pi_i(0, t)$ are obtained, $\pi_i(\tau, t)$ can be obtained by (20a).

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AUTHORS

Dr. Yashwant K. Malaiya; Department of Computer Science; School of Advanced Technology; State University of New York; Binghamton, New York 13901 USA.

Yashwant Malaiya (S'76, S'78, M'78) was born in 1951 in Saugor, India. He received his BS in 1969, MS in Physics in 1971, both from the University of Saugor. He obtained the MS (Tech.) in electronics from BITS, Pilani in 1974 and his PhD in electrical engineering from Utah State University in 1978. He is an assistant professor in the Department of Computer Science. His interests are intermittent faults, reliability analysis, fault-tolerance in multiprocessor systems, and design automation.

Prof. Stephen Y.H. Su; Department of Computer Science; School of Advanced Technology; State University of New York; Binghamton, New York 13901 USA.

Stephen Su (S'65, M'67, SM'77) was born in 1938 in Fukien, China. He received his BS in electrical engineering from the National Taiwan University in 1960, and the MS in electrical engineering and his PhD degree in computer engineering both from the University of Wisconsin, Madison in 1963 and 1967, respectively. He has been a chairman of two

international symposia and is a senior member of IEEE. He is a professor in the Department of Computer Science. His interests are intermittent faults, short-circuit and delay faults, design automation, and fault-tolerant design.

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